

Backstepping boundary control of Burgers' equation with actuator dynamics [☆]

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Abstract

In this paper, we propose a backstepping boundary control law for Burgers' equation with actuator dynamics. While the control law without actuator dynamics depends only on the signals $u(0, t)$ and $u(1, t)$, the backstepping control also depends on $u_x(0, t)$, $u_x(1, t)$, $u_{xx}(0, t)$ and $u_{xx}(1, t)$, making the regularity of the control inputs the key technical issue of the paper. With elaborate Lyapunov analysis, we prove that all these signals are sufficiently regular and the closed-loop system, including the boundary dynamics, is globally H^3 stable and well posed. © 2000 Elsevier Science B.V. All rights reserved.

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1. Introduction

In this paper, we develop a backstepping boundary feedback control for Burgers' equation

$$u_t - \varepsilon u_{xx} + uu_x = 0, \quad 0 < x < 1, \quad t > 0, \quad (1.1)$$

$$u_{xt}|_{x=0} = \varphi_0, \quad t > 0, \quad (1.2)$$

$$u_{xt}|_{x=1} = \varphi_1, \quad t > 0, \quad (1.3)$$

$$u|_{t=0} = u^0, \quad 0 < x < 1, \quad (1.4)$$

where $\varepsilon > 0$ is a viscosity parameter. In ODEs (1.2) and (1.3), φ_0 and φ_1 are control inputs. The function $u^0 = u^0(x)$ is an initial state in an appropriate function space. Since the above system is composed of a PDE

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(the viscous Burgers equation) and two ODEs, it is often referred to as a hybrid system in the literature (see, e.g., [16]).

The integrators separating φ_0 from $u_x|_{x=0}$ and φ_1 from $u_x|_{x=1}$ can be regarded as a part of actuator dynamics which prevent direct actuation via boundary values of u_x . The focus on Burgers' equation is much less for its physical relevance than for its structural properties that allow to address nontrivial nonlinear issues in a notationally simple setting, as a start towards future more practical nonlinear PDE stabilization problems.

To see that control is indeed needed for system (1.1)–(1.4), consider the uncontrolled system (i.e., $\varphi_0 = \varphi_1 = 0$)

$$u_t - \varepsilon u_{xx} + uu_x = 0, \quad 0 < x < 1, \quad t > 0, \quad (1.5)$$

$$u_{xt}|_{x=0} = 0, \quad t > 0, \quad (1.6)$$

$$u_{xt}|_{x=1} = 0, \quad t > 0, \quad (1.7)$$

$$u|_{t=0} = u^0, \quad 0 < x < 1 \quad (1.8)$$

with special initial conditions

$$u_x|_{x=0,t=0} = u_x|_{x=1,t=0} = \frac{2\varepsilon}{\cos^2(1/2)}. \quad (1.9)$$

To find the equilibrium points of (1.5)–(1.9), let us solve the steady equation of (1.5)–(1.9)

$$- \varepsilon v_{xx} + vv_x = 0, \quad 0 < x < 1, \quad (1.10)$$

$$v_x(0) = v_x(1) = \frac{2\varepsilon}{\cos^2(1/2)}. \quad (1.11)$$

Eqs. (1.10) and (1.11) can be written as

$$(\varepsilon v_x - \frac{1}{2}v^2)_x = 0$$

and then

$$\varepsilon v_x - \frac{1}{2}v^2 = C,$$

where C is a constant. Taking $C = 4\varepsilon^2$, we obtain one of the solutions of (1.10) and (1.11)

$$v(x) = 2\varepsilon \tan\left(x + \frac{\pi - 1}{2}\right), \quad (1.12)$$

which blows up at $x = \frac{1}{2}$.

The problem of control of Burgers' equation has received extensive attention recently [2–5,7,8,10,15]. In the present paper, we propose a backstepping control building upon the design

$$u_x(0, t) = k[u(0, t)^3 + u(0, t)], \quad (1.13)$$

$$u_x(1, t) = -k[u(1, t)^3 + u(1, t)], \quad (1.14)$$

in [10] where global boundary feedback stabilization was achieved for the positive constant k large enough. While controls (1.13) and (1.14) use only $u(0, t)$ and $u(1, t)$ for feedback, the backstepping controls applied through integrators (1.2) and (1.3) will employ also $u_x(0, t)$, $u_x(1, t)$, $u_{xx}(0, t)$ and $u_{xx}(1, t)$. To establish the regularity of those variables, which is the main subject of the analysis in this paper, we will show $H^3(0, 1)$ regularity of u .

To our knowledge, Coron and d’Andrea-Novel [6] were the first to apply backstepping to a PDE system. Common features our work has with theirs are that

- both handle actuator dynamics with backstepping,
- both show regularity of control inputs.

The distinguishing features of our work are:

- while [6] deals with a hyperbolic PDE, our paper is the first dealing with a parabolic one,
- while both [6] and our design employ scalar control inputs, theirs is applied in a distributed fashion, whereas ours acts only from the boundary.²

The rest of the paper is organized as follows. We design a backstepping boundary control and present our main results in Section 2. By using the Lyapunov method, we prove our main results in Section 3.

Notation: $H^s(0, 1)$ denotes the usual Sobolev space (see [1,13]) for any $s \in \mathbb{R}$. For $s \geq 0$, $H_0^s(0, 1)$ denotes the completion of $C_0^\infty(0, 1)$ in $H^s(0, 1)$, where $C_0^\infty(0, 1)$ denotes the space of all infinitely differentiable functions on $(0, 1)$ with compact support in $(0, 1)$. We use the following H^1 norm of $H^1(0, 1)$:

$$\|u\|_{H^1} = \left(u(0)^2 + \int_0^1 u_x^2 dx \right)^{1/2}, \quad u \in H^1(0, 1),$$

which is equivalent to the usual one. The norm on $L^2(0, 1)$ is denoted by $\|\cdot\|$. It is easy to see that

$$\|u\|^2 \leq 2\|u\|_{H^1}^2. \tag{1.15}$$

Let X be a Banach space and $T > 0$. We denote by $C^n([0, T]; X)$ the space of n times continuously differentiable functions defined on $[0, T]$ with values in X , and write $C([0, T]; X)$ for $C^0([0, T]; X)$. In what follows, for simplicity, we omit the indication of the varying range of x and t in equations and we understand that x varies from 0 to 1 and t from 0 to ∞ .

2. Main result

For notational convenience, in what follows, we denote

$$w_0 = u|_{x=0}, \quad w_1 = u|_{x=1}. \tag{2.1}$$

In order to use the backstepping method, we introduce the errors z_0, z_1 of control (1.13) and (1.14) as follows:

$$z_0 = u_x|_{x=0} - k(w_0 + w_0^3), \tag{2.2}$$

$$z_1 = u_x|_{x=1} + k(w_1 + w_1^3), \tag{2.3}$$

where k is a positive constant. With (1.1)–(1.3), we have

$$\dot{z}_0 = \varphi_0 - k(1 + 3w_0^2)(\varepsilon u_{xx} - uu_x)|_{x=0}, \tag{2.4}$$

$$\dot{z}_1 = \varphi_1 + k(1 + 3w_1^2)(\varepsilon u_{xx} - uu_x)|_{x=1}. \tag{2.5}$$

To achieve that the errors z_0 and z_1 decay exponentially, that is,

$$\dot{z}_0 = -\alpha z_0, \tag{2.6}$$

$$\dot{z}_1 = -\alpha z_1 \tag{2.7}$$

² Additionally, by nature of the two problems, [6] shows stability only in terms of basic energy, whereas we show stability in higher-order norms.

(where $\alpha > 0$), we choose the controls

$$\begin{aligned}\varphi_0 &= -\alpha z_0 + k(1 + 3w_0^2)(\varepsilon u_{xx} - uu_x)|_{x=0} \\ &= -\alpha(u_x|_{x=0} - k(w_0 + w_0^3)) + k(1 + 3w_0^2)(\varepsilon u_{xx} - uu_x)|_{x=0},\end{aligned}\quad (2.8)$$

$$\begin{aligned}\varphi_1 &= -\alpha z_1 - k(1 + 3w_1^2)(\varepsilon u_{xx} - uu_x)|_{x=1} \\ &= -\alpha(u_x|_{x=1} + k(w_1 + w_1^3)) - k(1 + 3w_1^2)(\varepsilon u_{xx} - uu_x)|_{x=1}.\end{aligned}\quad (2.9)$$

Now the closed-loop system (1.1), (2.2)–(2.3) and (2.6)–(2.9) can be summarized as

$$u_t - \varepsilon u_{xx} + uu_x = 0, \quad (2.10)$$

$$u_x|_{x=0} = k(w_0 + w_0^3) + z_0, \quad (2.11)$$

$$u_x|_{x=1} = -k(w_1 + w_1^3) + z_1, \quad (2.12)$$

$$\dot{z}_0 = -\alpha z_0, \quad (2.13)$$

$$\dot{z}_1 = -\alpha z_1 \quad (2.14)$$

which can be viewed as system (1.1), (1.13) and (1.14) perturbed by the exponentially decaying z_0 and z_1 . Hence, the *closed-loop* system differs from those previously analyzed in [2,14] in a very minor way. The novelty here is that the *control*

$$\varphi_0 = -\alpha z_0 + k(1 + 3w_0^2)(\varepsilon u_{xx}|_{x=0} - w_0(k(w_0 + w_0^3) + z_0)), \quad (2.15)$$

$$\varphi_1 = -\alpha z_1 - k(1 + 3w_1^2)(\varepsilon u_{xx}|_{x=1} - w_1(-k(w_1 + w_1^3) + z_1)), \quad (2.16)$$

besides employing the signals w_0, w_1, z_0, z_1 , also employs $u_{xx}|_{x=0}$ and $u_{xx}|_{x=1}$. To show that the latter signals, and hence the control, are sufficiently regular, we will require u^0 to be in $H^3(0, 1)$ and will need to perform substantial additional analysis.

Note that, in general, exponentially decaying perturbations can cause finite escape time phenomena in nonlinear systems [11], the prevention of which is ensured by employing the backstepping approach which forces error variables like z_0 and z_1 to decay faster than exponential, if needed. As we shall see, in the present problem, exponential decay of z_0 and z_1 will be sufficient to maintain global stability, which is a consequence of a specific way they enter Burgers' equation.

To show that control (2.8) and (2.9) indeed exponentially stabilizes system (1.1)–(1.4) at least in L^2 norm, we define the Lyapunov function

$$V = E + \frac{6\varepsilon^2}{\alpha} (z_0^2 + z_1^2), \quad (2.17)$$

where the energy function E is defined by

$$E = \int_0^1 u^2 dx. \quad (2.18)$$

Using (1.1)–(1.3) and integrating by parts, we obtain

$$\begin{aligned}\dot{E} &= 2 \int_0^1 u(\varepsilon u_{xx} - uu_x) dx \\ &= 2\varepsilon w_1[z_1 - k(w_1 + w_1^3)] - 2\varepsilon w_0[z_0 + k(w_0 + w_0^3)] - 2\varepsilon \int_0^1 u_x^2 dx - \frac{2}{3}(w_1^3 - w_0^3) \\ &\leq \left(\frac{1}{3} + \frac{1}{6} - 2\varepsilon k\right)(w_0^2 + w_0^4 + w_1^2 + w_1^4) + 6\varepsilon^2(z_0^2 + z_1^2) - 2\varepsilon \int_0^1 u_x^2 dx.\end{aligned}\quad (2.19)$$

Taking $k > 1/4\epsilon$ and denoting

$$E_1 = \int_0^1 u_x^2 dx + \frac{k}{2} (w_0^4 + 2w_0^2 + w_1^4 + 2w_1^2), \tag{2.20}$$

we obtain

$$\begin{aligned} \dot{E} &\leq \left(\frac{1}{2} - 2\epsilon k\right) (w_0^2 + w_0^4 + w_1^2 + w_1^4) + 6\epsilon^2(z_0^2 + z_1^2) - 2\epsilon \int_0^1 u_x^2 dx \\ &\leq 6\epsilon^2(z_0^2 + z_1^2) - \min\{2\epsilon k - 1/2, 2\epsilon\} \left(\int_0^1 u_x^2 dx + w_0^2 + w_0^4 + w_1^2 + w_1^4\right) \\ &= 6\epsilon^2(z_0^2 + z_1^2) - \min\{2\epsilon k - 1/2, 2\epsilon\} \left(\int_0^1 u_x^2 dx + \frac{k}{2k}(2w_0^2 + 2w_0^4 + 2w_1^2 + 2w_1^4)\right) \\ &\leq 6\epsilon^2(z_0^2 + z_1^2) - \min\{2\epsilon k - 1/2, 2\epsilon\} \left(\int_0^1 u_x^2 dx + \frac{k}{2k}(2w_0^2 + w_0^4 + 2w_1^2 + w_1^4)\right) \\ &\quad \left(\text{note that } \frac{1}{1+k} < 1 \text{ and } \frac{1}{1+k} < \frac{1}{k}\right) \\ &\leq 6\epsilon^2(z_0^2 + z_1^2) - \frac{1}{1+k} \min\{2\epsilon k - 1/2, 2\epsilon\} \left(\int_0^1 u_x^2 dx + \frac{k}{2}(2w_0^2 + w_0^4 + 2w_1^2 + w_1^4)\right) \\ &= 6\epsilon^2(z_0^2 + z_1^2) - \beta E_1, \end{aligned} \tag{2.21}$$

where

$$\beta = \frac{1}{1+k} \min\{2\epsilon k - 1/2, 2\epsilon\} > 0. \tag{2.22}$$

It therefore follows from (2.6) and (2.7) that

$$\begin{aligned} \dot{V} &= \dot{E} + \frac{12\epsilon^2}{\alpha} (z_0\dot{z}_0 + z_1\dot{z}_1) \\ &\leq 6\epsilon^2(z_0^2 + z_1^2) - \beta E_1 - 12\epsilon^2(z_0^2 + z_1^2) \\ &\leq -\gamma(E_1 + z_0^2 + z_1^2), \end{aligned} \tag{2.23}$$

where

$$\gamma = \min\{\beta, 6\epsilon^2\}. \tag{2.24}$$

Inequality (2.23) shows the L^2 exponential stability. Further, the closed-loop system

$$u_t - \epsilon u_{xx} + uu_x = 0, \tag{2.25}$$

$$u_{xt}|_{x=0} = [-\alpha(u_x - k(u + u^3)) + k(1 + 3u^2)(\epsilon u_{xx} - uu_x)]|_{x=0}, \tag{2.26}$$

$$u_{xt}|_{x=1} = [-\alpha(u_x + k(u + u^3)) - k(1 + 3u^2)(\epsilon u_{xx} - uu_x)]|_{x=1}, \tag{2.27}$$

$$u|_{t=0} = u^0 \tag{2.28}$$

is H^3 stable and well posed, as stated in the following theorem.

Theorem 2.1. *Suppose that the initial data $u^0 \in H^3(0, 1)$ and $k > 1/(4\epsilon)$. Let z_0 and z_1 be defined by (2.2) and (2.3), respectively. Then problem (2.25)–(2.28) has a unique global classical solution u with*

$$u \in C([0, \infty); H^3(0, 1)) \cap C^1([0, \infty); H^1(0, 1)).$$

Moreover, the solution satisfies the following L^2 , H^1 and H^3 stability estimates

$$\|u(t)\|^2 + \frac{6\varepsilon^2}{\alpha} [z_0(t)^2 + z_1(t)^2] \leq \left(\|u^0\|^2 + \frac{6\varepsilon^2}{\alpha} [z_0(0)^2 + z_1(0)^2] \right) e^{-2\omega_0 t}, \quad (2.29)$$

$$\begin{aligned} \|u(t)\|_{H^1}^2 + u_x(0,t)^2 + u_x(1,t)^2 &\leq C(\|u^0\|_{H^1}^2 + u_x^0(0)^2 + u_x^0(1)^2) \\ &\quad \times \exp(C(\|u^0\|_{H^1}^2 + u_x^0(0)^2 + u_x^0(1)^2)^3) e^{-\omega_1 t}, \end{aligned} \quad (2.30)$$

$$\|u(t)\|_{H^3}^2 \leq C(\|u^0\|_{H^3}^2) \exp(C\|u^0\|_{H^3}^6) e^{-\omega_2 t} \quad (2.31)$$

for $t \geq 0$, where

$$0 < \omega_0 = \frac{k\gamma}{4(1+k)\max\{1, 6\varepsilon^2/\alpha\}} < \varepsilon/2, \quad (2.32)$$

$$0 < \omega_1 < \min \left\{ \omega_0, \frac{k^2\alpha\gamma}{(3k+8)(4\alpha(1+k) + 6k\varepsilon^2)} \right\}, \quad (2.33)$$

$$\omega_2 = \min\{\omega_1, 2\varepsilon/5, \varepsilon k/5\} \quad (2.34)$$

and γ is defined by (2.24) and $C = C(\varepsilon, \alpha, k)$ is a positive constant.

Remark 2.1. In the case where $u^0 \in H^3(0, 1)$, it can be seen from (2.15), (2.16) and (2.31) that the control inputs φ_0 and φ_1 are bounded and converge to zero exponentially as $t \rightarrow \infty$.

3. Proof

In this section we prove our main result. By applying Theorem 7.4 of [12, p. 491], we show that problem (2.10)–(2.14), which is equivalent to (2.25)–(2.28), is well posed. The stability estimates (2.30) and (2.31) are proved with the Lyapunov method.

Proof of Theorem 2.1. *Step 1: Well-posedness.* Solving ordinary differential equations (2.26) and (2.27), we obtain

$$u_x|_{x=0} = z_0^0 e^{-\alpha t} + k(w_0 + w_0^3), \quad (3.1)$$

$$u_x|_{x=1} = z_1^0 e^{-\alpha t} - k(w_1 + w_1^3), \quad (3.2)$$

where

$$z_0^0 = u_x^0(0) - k[u^0(0) + u^0(0)^3], \quad (3.3)$$

$$z_1^0 = u_x^0(1) + k[u^0(1) + u^0(1)^3]. \quad (3.4)$$

Consequently, in the class of classical solutions, problem (2.25)–(2.28) is equivalent to

$$u_t - \varepsilon u_{xx} + uu_x = 0, \quad (3.5)$$

$$u_x|_{x=0} = z_0^0 e^{-\alpha t} + k(w_0 + w_0^3), \quad (3.6)$$

$$u_x|_{x=1} = z_1^0 e^{-\alpha t} - k(w_1 + w_1^3), \quad (3.7)$$

$$u|_{t=0} = u^0. \quad (3.8)$$

Set

$$a(x, t) = x(x-1)^2 z_0^0 e^{-\alpha t} + x^2(x-1) z_1^0 e^{-\alpha t}, \quad (3.9)$$

$$v(x, t) = u(x, t) - a(x, t). \tag{3.10}$$

Then we have

$$\begin{aligned} v_x|_{x=0} &= u_x|_{x=0} - a_x|_{x=0} \\ &= z_0^0 e^{-\alpha t} + k(w_0 + w_0^3) - z_0^0 e^{-\alpha t} \\ &\quad (\text{note that } v(0, t) = u(0, t) = w_0) \\ &= k(v(0, t) + v(0, t)^3) \end{aligned} \tag{3.11}$$

and likewise,

$$v_x|_{x=1} = -k(v(1, t) + v(1, t)^3). \tag{3.12}$$

Hence, problem (3.5)–(3.8) becomes

$$v_t - \varepsilon v_{xx} + (v + a)(v_x + a_x) + a_t - \varepsilon a_{xx} = 0, \tag{3.13}$$

$$v_x|_{x=0} = k(v(0, t) + v(0, t)^3), \tag{3.14}$$

$$v_x|_{x=1} = -(v(1, t) + v(1, t)^3), \tag{3.15}$$

$$v|_{t=0} = u^0 - a(x, 0). \tag{3.16}$$

We can readily verify that problem (3.13)–(3.16) satisfies all conditions of Theorem 7.4 of [12, p. 491] with any positive constant β there (especially note that condition (7.36) on p. 487 of [12] is satisfied). Therefore we conclude from this theorem that problem (3.13)–(3.16) has a unique classical solution v in the class $H^{2+\beta, 1+\beta/2}(\bar{\Omega} \times [0, T])$ for suitable $T > 0$. Consequently, problem (3.5)–(3.8) has a unique classical solution $u = v + a$ in the class $H^{2+\beta, 1+\beta/2}(\bar{\Omega} \times [0, T])$. Global solutions are obtained by establishing a priori estimates below.

Step 2: L^2 -estimate. In what follows, we denote by $C = C(\varepsilon, \alpha, k)$ a generic positive constant that may vary from line to line.

Since

$$V \leq \frac{2(1+k)}{k} \max\{1, 6\varepsilon^2/\alpha\}(E_1 + z_0^2 + z_1^2), \tag{3.17}$$

we deduce from (2.23) that

$$\dot{V} \leq -\frac{k\gamma}{2(1+k)\max\{1, 6\varepsilon^2/\alpha\}} V, \tag{3.18}$$

which implies that

$$V \leq V(0)e^{-2\omega_0 t}, \quad t \geq 0, \tag{3.19}$$

where γ and ω_0 are defined by (2.24) and (2.32), respectively. This proves (2.29).

Step 3: H^1 -estimate. On the other hand, integrating by parts and using Eqs. (3.5)–(3.7), we obtain

$$\begin{aligned} \dot{E}_1 &= 2 \int_0^1 u_x u_{xt} \, dx + 2k[(w_1^3 + w_1)\dot{w}_1 + (w_0^3 + w_0)\dot{w}_0] \\ &= 2\dot{w}_1 z_1 - 2\dot{w}_0 z_0 - 2 \int_0^1 u_{xx}(\varepsilon u_{xx} - uu_x) \, dx \\ &\quad \times \left(\text{note that } \max_{0 \leq x \leq 1} u(x)^2 \leq \frac{2(1+k)}{k} E_1 \text{ and } 2 \int_0^1 u_{xx} uu_x \, dx \leq \varepsilon \int_0^1 u_{xx}^2 + \frac{1}{\varepsilon} \max_{0 \leq x \leq 1} u(x)^2 \int_0^1 u_x^2 \, dx \right) \\ &\leq 2z_1 \dot{w}_1 - 2z_0 \dot{w}_0 - \varepsilon \int_0^1 u_{xx}^2 \, dx + \frac{2(1+k)}{\varepsilon k} E_1^2 \\ &\leq 2 \frac{d}{dt}(w_1 z_1) + 2\alpha z_1 w_1 - 2 \frac{d}{dt}(w_0 z_0) - 2\alpha z_0 w_0 + \frac{2(1+k)}{\varepsilon k} E_1^2. \end{aligned} \tag{3.20}$$

Adding $d/dt(\frac{2}{k}z_0^2)$ and $d/dt(\frac{2}{k}z_1^2)$ to both sides, we obtain

$$\begin{aligned} & \frac{d}{dt} \left(E_1 - \frac{k}{2} (w_0^2 + w_1^2) + \frac{2}{k} \left(\frac{k}{2} w_0 + z_0 \right)^2 + \frac{2}{k} \left(\frac{k}{2} w_1 - z_1 \right)^2 \right) \\ & \leq 2\alpha z_1 w_1 - 2\alpha z_0 w_0 - \frac{4\alpha}{k} (z_0^2 + z_1^2) + \frac{2(1+k)}{\varepsilon k} E_1^2 \\ & \leq C(w_0^2 + w_1^2) + \frac{2(1+k)}{\varepsilon k} E_1^2. \end{aligned} \quad (3.21)$$

Denoting

$$\begin{aligned} J &= E_1 - \frac{k}{2} (w_0^2 + w_1^2) + \frac{2}{k} \left(\frac{k}{2} w_0 + z_0 \right)^2 + \frac{2}{k} \left(\frac{k}{2} w_1 - z_1 \right)^2 \\ &= \int_0^1 u_x^2 + \frac{k}{2} (w_0^2 + w_0^4 + w_1^2 + w_1^4) \\ & \quad + \frac{2}{k} \left[\left(u_x|_{x=0} - \frac{k}{2} w_0 - k w_0^3 \right)^2 + \left(u_x|_{x=1} + \frac{k}{2} w_1 + k w_1^3 \right)^2 \right], \end{aligned} \quad (3.22)$$

we get

$$J \leq C(w_0^2 + w_1^2) + \frac{2(1+k)}{\varepsilon k} E_1^2. \quad (3.23)$$

It is clear that

$$\frac{1}{2} E_1 \leq J \leq \frac{3}{2} E_1 + \frac{4}{k} (z_0^2 + z_1^2) \leq \frac{3k+8}{2k} (E_1 + z_0^2 + z_1^2). \quad (3.24)$$

For $M > 0$ yet to be determined, we define the Lyapunov function G as

$$G = MV + \ln \left(1 + J + \frac{6\varepsilon^2}{\alpha} (z_0^2 + z_1^2) \right), \quad (3.25)$$

which, in view of (3.22) and (3.24), is positive definite and radially unbounded in the argument of $\|u\|_{H^1}$, $u_x(0)$ and $u_x(1)$. It is also easy to see that

$$\begin{aligned} G &\leq MV + J + \frac{6\varepsilon^2}{\alpha} (z_0^2 + z_1^2) \\ &\leq \frac{2M(1+k)}{k} E_1 + \frac{6M\varepsilon^2}{\alpha} (z_0^2 + z_1^2) + J + \frac{6\varepsilon^2}{\alpha} (z_0^2 + z_1^2) \\ &\leq \frac{4M(1+k)}{k} J + \frac{6M\varepsilon^2}{\alpha} (z_0^2 + z_1^2) + J + \frac{6\varepsilon^2}{\alpha} (z_0^2 + z_1^2) \\ &\leq \frac{4M(1+k) + k}{k} J + \frac{6\varepsilon^2(M+1)}{\alpha} (z_0^2 + z_1^2) \\ &\leq \frac{\alpha(4M(1+k) + k) + 6k\varepsilon^2(M+1)}{k\alpha} (J + z_0^2 + z_1^2). \end{aligned} \quad (3.26)$$

It then follows from (2.23), (3.23) and (3.24) that for $M > 16(1+k)(3k+8)/2\varepsilon\gamma k^2$

$$\begin{aligned}
 \dot{G} &= M\dot{V} + \frac{J + (12\varepsilon^2/\alpha)(z_0\dot{z}_0 + z_1\dot{z}_1)}{1 + J + (6\varepsilon^2/\alpha)(z_0^2 + z_1^2)} \\
 &\quad (\text{use (3.23) and note that } z_0\dot{z}_0 + z_1\dot{z}_1 = -\alpha(z_0^2 + z_1^2) \text{ is negative)} \\
 &\leq M\dot{V} + \frac{C(w_0^2 + w_1^2) + [2(1+k)/\varepsilon k]E_1^2}{1 + J + (6\varepsilon^2/\alpha)(z_0^2 + z_1^2)} \\
 &\quad (\text{use (2.23) and the first inequality of (3.24)}) \\
 &\leq -M\gamma(E_1 + z_0^2 + z_1^2) + \frac{8(1+k)}{\varepsilon k} J + C(w_0^2 + w_1^2) \\
 &\leq -\frac{M\gamma}{2}(E_1 + z_0^2 + z_1^2) - \frac{M\gamma}{2}(z_0^2 + z_1^2) + \frac{8(1+k)}{\varepsilon k} J + \left(C - \frac{Mk\gamma}{2}\right)(w_0^2 + w_1^2) \\
 &\leq -\frac{2kM\gamma}{2(3k+8)} J - \frac{M\gamma}{2}(z_0^2 + z_1^2) + \frac{8(1+k)}{\varepsilon k} J + \left(C - \frac{Mk\gamma}{2}\right)(w_0^2 + w_1^2) \\
 &\leq -\left(\frac{2kM\gamma}{2(3k+8)} - \frac{8(1+k)}{\varepsilon k}\right)(J + z_0^2 + z_1^2) + \left(C - \frac{Mk\gamma}{2}\right)(w_0^2 + w_1^2) \\
 &\leq -\frac{k\alpha}{\alpha(4M(1+k) + k) + 6k\varepsilon^2(M+1)} \left(\frac{2kM\gamma}{2(3k+8)} - \frac{8(1+k)}{\varepsilon k}\right) G \\
 &\quad + \left(C - \frac{Mk\gamma}{2}\right)(w_0^2 + w_1^2). \tag{3.27}
 \end{aligned}$$

Since the function

$$f(M) = \frac{k\alpha}{\alpha(4M(1+k) + k) + 6k\varepsilon^2(M+1)} \left(\frac{2kM\gamma}{2(3k+8)} - \frac{8(1+k)}{\varepsilon k}\right)$$

is increasing and

$$\lim_{M \rightarrow \infty} f(M) = \frac{k^2\alpha\gamma}{(3k+8)(4\alpha(1+k) + 6k\varepsilon^2)},$$

we can find an M_0 sufficiently large such that $\omega_1 < f(M_0)$ and $C - M_0k\gamma/2 < 0$, where ω_1 is defined by (2.33). Consequently, we obtain

$$\dot{G} \leq -\omega_1 G. \tag{3.28}$$

Hence, we have

$$G \leq G(0)e^{-\omega_1 t}, \quad t \geq 0. \tag{3.29}$$

It therefore follows that

$$1 + J + \frac{6\varepsilon^2}{\alpha}(z_0^2 + z_1^2) \leq \exp(G(0)e^{-\omega_1 t}),$$

and then

$$\begin{aligned}
 J + \frac{6\varepsilon^2}{\alpha}(z_0^2 + z_1^2) &\leq \exp(G(0)e^{-\omega_1 t}) - 1 \\
 &\leq CG(0)\exp(G(0))e^{-\omega_1 t} \\
 &\leq C(\|u^0\|_{H^1}^2 + \|u^0\|_{H^1}^6 + u_x^0(0)^2 + u_x^0(1)^2) \\
 &\quad \times \exp(C(\|u^0\|_{H^1}^2 + \|u^0\|_{H^1}^6 + u_x^0(0)^2 + u_x^0(1)^2))e^{-\omega_1 t}, \quad t \geq 0, \tag{3.30}
 \end{aligned}$$

which implies (2.30) since

$$\begin{aligned}
 & \|u^0\|_{H^1}^2 + \|u^0\|_{H^1}^6 + u_x^0(0)^2 + u_x^0(1)^2 \\
 & \leq \|u^0\|_{H^1}^2 + u_x^0(0)^2 + u_x^0(1)^2 + (\|u^0\|_{H^1}^2 + u_x^0(0)^2 + u_x^0(1)^2)^3 \\
 & \leq C(\|u^0\|_{H^1}^2 + u_x^0(0)^2 + u_x^0(1)^2) \exp(\|u^0\|_{H^1}^2 + u_x^0(0)^2 + u_x^0(1)^2) \\
 & \leq eC(\|u^0\|_{H^1}^2 + u_x^0(0)^2 + u_x^0(1)^2) \exp((\|u^0\|_{H^1}^2 + u_x^0(0)^2 + u_x^0(1)^2)^3)
 \end{aligned} \tag{3.31}$$

and

$$\begin{aligned}
 & \exp(C(\|u^0\|_{H^1}^2 + \|u^0\|_{H^1}^6 + u_x^0(0)^2 + u_x^0(1)^2)) \\
 & \leq \exp(C(1 + (\|u^0\|_{H^1}^2 + u_x^0(0)^2 + u_x^0(1)^2)^3)) \\
 & = e^C \exp(C(\|u^0\|_{H^1}^2 + u_x^0(0)^2 + u_x^0(1)^2)^3).
 \end{aligned} \tag{3.32}$$

Step 4: H^2 -estimate. In order to obtain the H^2 -estimate, we estimate the L^2 norm of u_t . Integrating by parts and using Eqs. (3.5)–(3.7), we obtain

$$\begin{aligned}
 \frac{d}{dt} \int_0^1 u_t^2 dx &= 2 \int_0^1 u_t(\varepsilon u_{txx} - u_t u_x - u u_{tx}) dx \\
 &\leq -2\varepsilon k[\dot{w}_0^2(3w_0^2 + 1) + \dot{w}_1^2(3w_1^2 + 1)] + 2\varepsilon \dot{w}_1 \dot{z}_1 - 2\varepsilon \dot{w}_0 \dot{z}_0 - 2\varepsilon \|u_{xt}\|^2 + C\|u_t\| \|u_t\|_{H^1} \|u\|_{H^1} \\
 &= -2\varepsilon k[\dot{w}_0^2(3w_0^2 + 1) + \dot{w}_1^2(3w_1^2 + 1)] \\
 &\quad - 2\varepsilon \alpha \dot{w}_1 z_1 + 2\varepsilon \alpha \dot{w}_0 z_0 \quad (\text{use (2.6) and (2.7)}) \\
 &\quad - 2\varepsilon \|u_{xt}\|^2 + C\|u_t\| \|u_t\|_{H^1} \|u\|_{H^1} \\
 &\quad (\text{note definition (2.34) of } \omega_2) \\
 &\leq -\frac{1}{2}\varepsilon k[\dot{w}_0^2(3w_0^2 + 1) + \dot{w}_1^2(3w_1^2 + 1)] + C(z_1^2 + z_0^2) - 5\omega_2 \|u_t\|_{H^1}^2 + C\|u_t\| \|u_t\|_{H^1} \|u\|_{H^1} \\
 &\leq -\frac{1}{2}\varepsilon k[\dot{w}_0^2(3w_0^2 + 1) + \dot{w}_1^2(3w_1^2 + 1)] + C(z_1^2 + z_0^2) - 4\omega_2 \|u_t\|_{H^1}^2 + C\|u_t\|^2 \|u\|_{H^1}^2 \\
 &\leq -\frac{1}{2}\varepsilon k[\dot{w}_0^2(3w_0^2 + 1) + \dot{w}_1^2(3w_1^2 + 1)]C(z_1^2 + z_0^2) - 2\omega_2 \|u_t\|^2 + C\|u_t\|^2 \|u\|_{H^1}^2.
 \end{aligned} \tag{3.33}$$

Multiplying (3.33) by $e^{2\omega_2 t}$ gives

$$\begin{aligned}
 \frac{d}{dt} (\|u_t\|^2 e^{2\omega_2 t}) &\leq C(z_1^2 + z_0^2) e^{2\omega_2 t} + C\|u\|_{H^1}^2 \|u_t\|^2 e^{2\omega_2 t} \\
 &\leq C(z_1^2 + z_0^2) e^{2\omega_2 t} + C\|u\|_{H^1}^2 \|u_t\|^2 e^{2\omega_2 t}.
 \end{aligned} \tag{3.34}$$

It, therefore, follows from (2.29) that

$$\|u_t\|^2 e^{2\omega_2 t} \leq \|u_t(0)\|^2 + C(\|u^0\|_{H^2}^2 + \|u^0\|_{H^2}^6) + C \int_0^t \|u\|_{H^1}^2 \|u_t\|^2 e^{2\omega_2 s} ds,$$

and then by (2.29) and Gronwall’s inequality we obtain

$$\|u_t\|^2 \leq C(\|u^0\|_{H^2}^2 + \|u^0\|_{H^2}^6) \exp(C(\|u^0\|_{H^2}^2 + \|u^0\|_{H^2}^6)) e^{-2\omega_2 t}, \quad t \geq 0. \tag{3.35}$$

Since

$$\|u_{xx}\|^2 \leq C(\|u_t\|^2 + \|u\|_{H^1}^4), \tag{3.36}$$

it follows from (3.30), (3.35) and (3.36) that

$$\|u_{xx}\|^2 \leq C(\|u^0\|_{H^2}^2 + \|u^0\|_{H^2}^{12}) \exp(C(\|u^0\|_{H^2}^2 + \|u^0\|_{H^2}^6)) e^{-2\omega_2 t}, \quad t \geq 0. \tag{3.37}$$

Step 5: H^3 -estimate. To prove that (2.31), we estimate u_{xt} . Integrating by parts and using (3.5)–(3.7), we obtain

$$\begin{aligned}
 \frac{d}{dt} \int_0^1 u_{xt}^2 dx &= 2u_{xt}u_{tt} \Big|_0^1 - 2 \int_0^1 u_{xxt}(\varepsilon u_{txx} - u_t u_x - uu_{xt}) dx \\
 &\leq -k \left[\frac{d}{dt}(\dot{w}_0^2)(3w_0^2 + 1) + \frac{d}{dt}(\dot{w}_1^2)(3w_1^2 + 1) \right] \\
 &\quad + 2\dot{w}_1 \dot{z}_1 - 2\dot{w}_0 \dot{z}_0 - 2\varepsilon \|u_{xxt}\|^2 + 8\|u_{xxt}\| \|u_t\|_{H^1} \|u\|_{H^1} \\
 &\leq -k \frac{d}{dt} [\dot{w}_0^2(3w_0^2 + 1) + \dot{w}_1^2(3w_1^2 + 1)] + 6k(\dot{w}_0^3 w_0 + \dot{w}_1^3 w_1) \\
 &\quad - 2\alpha z_1 \dot{w}_1 + 2\alpha z_0 \dot{w}_0 \quad (\text{use (2.6) and (2.7)}) \\
 &\quad + C\|u_t\|_{H^1}^2 \|u\|_{H^1}^2 \\
 &= -k \frac{d}{dt} [\dot{w}_0^2(3w_0^2 + 1) + \dot{w}_1^2(3w_1^2 + 1)] + 6k(\dot{w}_0^3 w_0 + \dot{w}_1^3 w_1) \\
 &\quad - 2\alpha \frac{d}{dt}(\dot{w}_1 z_1) + 2\alpha \frac{d}{dt}(\dot{w}_0 z_0) - 2z_1 \alpha^2 \dot{w}_1 + 2z_0 \alpha^2 \dot{w}_0 + C\|u_t\|_{H^1}^2 \|u\|_{H^1}^2. \tag{3.38}
 \end{aligned}$$

Adding $d/dt(\alpha^2 z_0^2/k)$ and $d/dt(\alpha^2 z_1^2/k)$ to both sides and setting

$$J_1 = \int_0^1 u_{xt}^2 dx + 3k(\dot{w}_0^2 w_0^2 + \dot{w}_1^2 w_1^2) + (\sqrt{k}\dot{w}_1 + z_1/\sqrt{k})^2 + (\sqrt{k}\dot{w}_0 - z_0/\sqrt{k})^2, \tag{3.39}$$

we obtain

$$\begin{aligned}
 \dot{J}_1 &\leq 6k(\dot{w}_0^3 w_0 + \dot{w}_1^3 w_1) + 2|z_1| \alpha^2 |\dot{w}_1| + 2|z_0| \alpha^2 |\dot{w}_0| - \frac{2\alpha^3}{k} (z_0^2 + z_1^2) + C\|u_t\|_{H^1}^2 \|u\|_{H^1}^2 \\
 &\leq \underbrace{6k(\dot{w}_0^3 w_0 + \dot{w}_1^3 w_1)}_{=I} + C(\dot{w}_0^2 + \dot{w}_1^2) + C\|u\|_{H^1}^2 J_1. \tag{3.40}
 \end{aligned}$$

We now estimate I . Since for any $x, y \in [0, 1]$

$$v(x)^2 = v(y)^2 + \int_y^x (v^2)_\xi d\xi \leq v(y)^2 + \|v\| \|v_x\|,$$

we have

$$\max_{0 \leq x \leq 1} v(x)^2 \leq \|v\|^2 + \|v\| \|v_x\|. \tag{3.41}$$

It, therefore, follows that

$$\begin{aligned}
 |\dot{w}_i^3 w_i| &\leq |\dot{w}_i w_i| (\|u_t\|^2 + \|u_t\| \|u_{xt}\|) \\
 &\leq \frac{1}{2}(\dot{w}_i^2 w_i^2 + \|u_t\|^4 + \|u_t\|^2 \|u_{xt}\|^2), \quad i = 0, 1 \tag{3.42}
 \end{aligned}$$

and then

$$I \leq C \sum_{i=0}^1 \dot{w}_i^2 w_i^2 + C\|u_t\|^4 + C\|u_t\|^2 \|u_{xt}\|^2. \tag{3.43}$$

Thus, we deduce from (3.40) that

$$\dot{J}_1 \leq C \sum_{i=0}^1 \dot{w}_i^2 (w_i^2 + 1) + C\|u_t\|^4 + C(\|u_t\|^2 + \|u\|_{H^1}^2) J_1. \tag{3.44}$$

On the other hand, multiplying (3.33) by $e^{\omega_2 t}$, one obtains

$$\begin{aligned} \frac{d}{dt} (\|u_t\|^2 e^{\omega_2 t}) + \frac{1}{2} \varepsilon k [\dot{w}_0^2 (3w_0^2 + 1) + \dot{w}_1^2 (3w_1^2 + 1)] e^{\omega_2 t} \\ \leq C(z_1^2 + z_0^2) e^{\omega_2 t} + C\|u\|_{H^1}^2 \|u_t\|^2 e^{\omega_2 t} \\ \leq C(z_1^2 + z_0^2) e^{\omega_2 t} + C\|u\|_{H^1}^2 \|u_t\|^2 e^{\omega_2 t}. \end{aligned} \quad (3.45)$$

Integrating from 0 to ∞ and using (3.30) and (3.35), one obtains

$$\int_0^\infty \sum_{i=0}^1 \dot{w}_i^2 (w_i^2 + 1) e^{\omega_2 t} dt \leq C(\|u^0\|_{H^2}^2 + \|u^0\|_{H^2}^6). \quad (3.46)$$

It, therefore, follows from Gronwall's inequality (see, e.g., [9, p. 63]) that

$$J_1 \leq C(\|u^0\|_{H^3}^2 + \|u^0\|_{H^3}^{12}) \exp(C(\|u^0\|_{H^2}^2 + \|u^0\|_{H^2}^6)) e^{-\omega_2 t}, \quad t \geq 0. \quad (3.47)$$

Since

$$\|u_{xxx}\|^2 \leq C(\|u_{xt}\|^2 + \|u\|_{H^2}^4) \quad (3.48)$$

(2.31) follows from (2.30), (3.37) and (3.47) with the use of (3.31) and (3.32). \square

4. Conclusion

In summary, we propose here a backstepping boundary control law for Burgers' equation with actuator dynamics, which depends not only on the signals $u(0, t)$ and $u(1, t)$, but also on $u_x(0, t)$, $u_x(1, t)$, $u_{xx}(0, t)$ and $u_{xx}(1, t)$. With elaborate Lyapunov analysis, and applying a result from the the classical theory of nonlinear partial differential equations of parabolic type, we prove that the controlled closed-loop system is globally H^3 stable and has a unique global classical solution.

As we pointed out in the introduction, system (1.1)–(1.3) is only an example of a more general class of boundary control problems (yet to be investigated) tractable by backstepping. To give a preview of possible future extensions, we offer the system

$$u_t - \varepsilon u_{xx} + uu_x = 0, \quad (4.1)$$

$$u_{xt}|_{x=0} = f_{01} \left(u|_{x=0}, u_x|_{x=0}, u_{xx}|_{x=0}, \int_0^1 u dx \right) + \zeta_0, \quad (4.2)$$

$$u_{xt}|_{x=1} = f_{11} \left(u|_{x=1}, u_x|_{x=1}, u_{xx}|_{x=1}, \int_0^1 u dx \right) + \zeta_1, \quad (4.3)$$

$$\dot{\zeta}_0 = f_{02} \left(u|_{x=0}, u_x|_{x=0}, u_{xx}|_{x=0}, \int_0^1 u dx, \zeta_0 \right) + \varphi_0, \quad (4.4)$$

$$\dot{\zeta}_1 = f_{12} \left(u|_{x=1}, u_x|_{x=1}, u_{xx}|_{x=1}, \int_0^1 u dx, \zeta_1 \right) + \varphi_1, \quad (4.5)$$

where f_{ij} are smooth functions and φ_0, φ_1 are controls. The variables ζ_0, ζ_1 serve as “virtual controls” to the $u|_{x=0}, u|_{x=1}$ equations, effectively separating the inputs ϕ_0, ϕ_1 from the Burgers equation but not one but two cascaded integrators (on each end of the $[0, 1]$ interval). After designing the control laws in two steps of backstepping, one would need $H^5(0, 1)$ estimates on u to show regularity of the control signals. Adding more integrators would be standard insofar as the backstepping design task is concerned, however, regularity analysis would require estimates of increasingly high order.

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